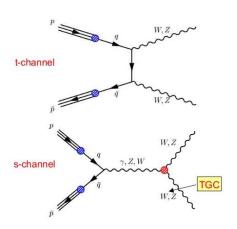
WW+WZ in $\ell\nu$ +jets

Viviana Cavaliere

Seminar

10/26/2010

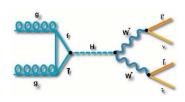
Why Diboson?

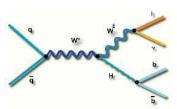


- Studies of electroweak (EW)
 vector boson production are an
 important aspect of the Tevatron
 physics program.
- Potential for new physics is manifest in:
 - Precision measurements of mass and EW parameters
 - Relationships between the masses of the W and Z
 - Increased cross sections or changes in kinematics
- Sensitive to new physics signatures

Motivation: Higgs searches

- ullet H o WW is the dominant decay mode for a high mass Higgs $(m_H>135~{
 m GeV}/c^2)$
 - Drives current exclusion limits
 - Direct diboson production is the single most important background
 - Importan to understand
- $WH \to \ell \nu b \bar{b}$ is a promising search mode for low mass Higgs $(m_H < 135~{
 m GeV}/c^2)$
 - Similar topology/final state to $WW/WZ \rightarrow \ell \nu q \bar{q}$
 - Similar challenges \rightarrow S/B WH 1.2% WW/WZ 2.9%



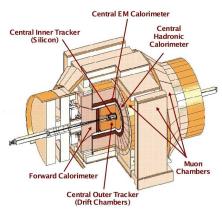


 ${\sf WW/WZ} {
ightarrow} \ \ell \nu q \bar{q}$ is a proving ground for Higgs search

CDF detector

- $\begin{tabular}{ll} \hline \bullet & {\rm Proton-antiproton \ collision \ at \ } \sqrt{s} = 1.96 \\ {\rm TeV} \\ \hline \end{tabular}$
- 36 bunches: crossing time = 396 ns
- $\ \, \textbf{P} \mbox{eak luminosity } 4\cdot 10^{32} \mbox{ cm}^{-2} \mbox{ s}^{-1} \label{eq:local_equation}$





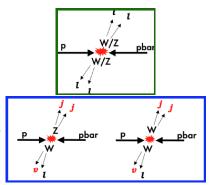
• About 8 fb^{-1} on tape

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• This analysis uses 4.3 fb^{-1}

Diboson final states

- WW, WZ and ZZ production has already been seen in leptonic final states:
 - Clean Yields but low BR
- Semi-leptonic modes suffer from large background:
 - WW, WZ, ZZ seen in \cancel{E}_T + jets mode
- Looking at $\ell\nu$ + jets final state:

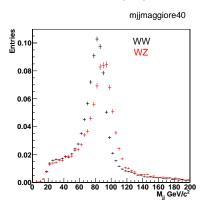


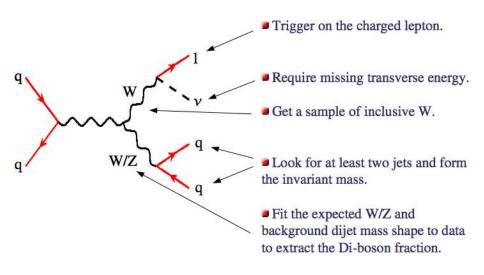
Signal:
$$\sigma(p\bar{p} \to WW/WZ) = 15.9 \pm 0.9 \text{ pb}$$

Background: $\sigma(p\bar{p} \to W + jets) = 2066 \text{ pb}$
 $\sigma(p\bar{p} \to Z + jets) = 187 \text{ pb}$

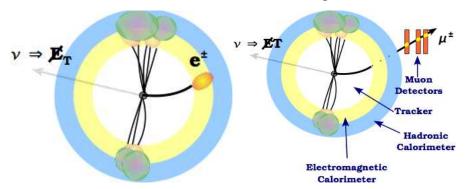
WW vs WZ

- We treat events from WW and WZ as indistinguishable signals
 - Largely due to insufficient dijet mass resolution: 10 GeV difference in mass
- Consider the relative selection efficiency for WW vs WZ:
 - WW (WZ) $\rightarrow \ell \nu j j$ branching fraction: \sim 28.5 (14.2)%
 - WW (WZ) $\rightarrow \ell \nu jj \ \sigma \cdot BR$: \sim 3.5 (0.5) pb

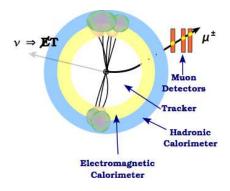




- One lepton (electron or muon) with $E_T/p_T > 20$ GeV and $|\eta| < 1.2$.
- Undetected neutrino manifests as an imbalance in transverse momentum: "missing" transverse energy ($\cancel{E}_T > 25$ GeV)
- ullet To reduce multijet backgrounds, we require $M_T^W >$ 30 GeV.



- Quark jets arising from $W/Z \rightarrow q\bar{q}$ decays are very energetic and relatively central
- Cluster energy in cones of $\Delta R <$ 0.4
- Calorimeter signature must be inconsistent with electron signatures
- \bullet $E_T > 20$ GeV and $|\eta| < 2.4$



Backgrounds

- $W \to \ell \nu + \text{jets } (l = e, \mu, \tau)$:
 - same signature as signal with a much higher cross section.
 - Almost 80% of the sample
 - Shape is taken from Alpgen MC
- $Z \rightarrow ll + \text{jets } (l = e, \mu, \tau)$:
 - ullet where one of the two leptons escapes detection and produces \mathcal{E}_T
 - shape taken from Alpgen MC.
- $t\bar{t}$ + single top: shape taken from Pythia MC.
- QCD Multijet :
 - e.g a three-jet event in which one jet passes all lepton cuts and, simultaneously, the energies are so badly measured that a large \mathcal{E}_T is reported.



QCD estimation

"AntiElectrons":

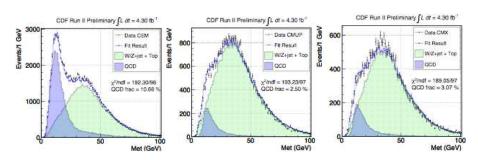
- Some non-kinematic cuts for the electron (EHAD/EM ...) are used to reject fake electrons.
- Model is constructed of events which fail at least two of the non-kinematic cuts but pass all the kinematic cuts of the electron.

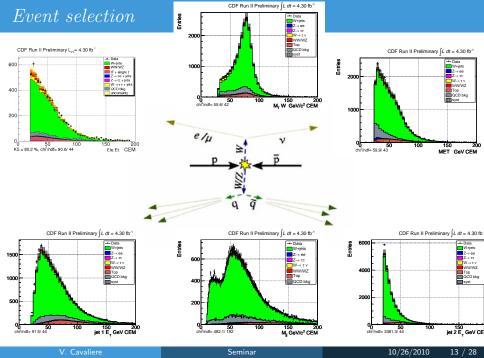
"Non isolated muons":

- Using non-isolated events, events which pass all selection criteria except the requirement of lepton isolation.
- Based on the rationale that non-isolated events are typically leptons contained in jets, and jets that contain energetic leptons are more likely to pass lepton identification cuts.

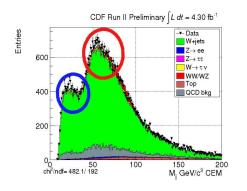
QCD estimation

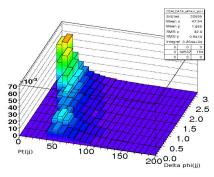
- QCD multi-jet events do not have neutrinos so met distribution is completely different from W+jets
- Extract the fraction of qcd and knowing all the others contributions can extract also the W+jets normalization.



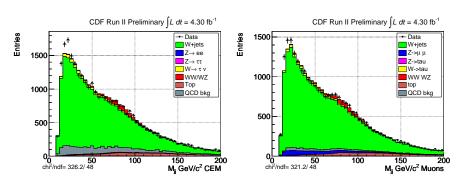


Di-jet mass





- The E_T threshold on the jets gives rise to two peaks:
 - At $m_{jj} \sim$ 20 GeV for almost collinear jets with $\Delta \phi \sim$ 0.5 where the invariant mass is minimum and the combined $P_{T,jj}$ is maximum
 - The second one is at $m_{jj} \sim$ 40 GeV, for back to back jets $(\Delta \phi \sim \pi)$, where the invariant mass is maximum.
- Decide to cut at $P_{T,jj} > 40 \text{ GeV/c} \rightarrow \text{loss of} \sim 40 \%$ of the signal



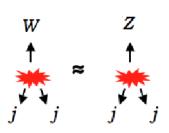
- Want to fit the m_{ij} distribution
- Need to be sure of the shape of every component

Checking the background shapes

- W+jets
- QCD backgrounds

$Checking\ W+jets\ shape$

- Use Z+jets data to check ewk shape
- Basid Idea: Similar kinematics
- Kinematics of Z+jets not necessarily the same of W +jets \rightarrow need to account for this

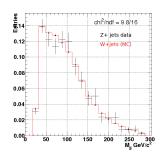


Method:

$$Z + jet(data) \approx \frac{Z + jets(MC)}{W + jets(MC)} \cdot W + jets(data)$$

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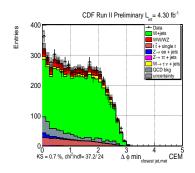


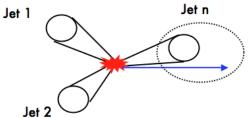
Method:

$$Z + jet(data) \approx \frac{Z + jets(MC)}{W + jets(MC)} \cdot W + jets(data)$$

Checking the QCD shape I

- ullet QCD o less understood background
- Need independent ways of checking the shape and rate
- Variable sensitive to qcd shape is $\Delta \phi(E_T \text{closest jet})$



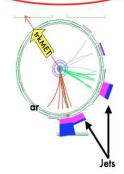


Mis-measured jet tends to align with MET → Δφ should help rejection

Checking the QCD II

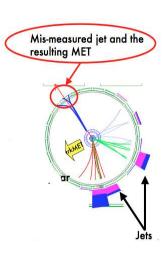
- E_T is normally calculated at calorimeter level
- Can also estimate it using tracks
- ullet True $ot\!\!\!E_T: \Delta\phi(
 ot\!\!\!E_T, {\sf trkmet})$ small
- Fake \mathscr{E}_T : $\Delta \phi(\mathscr{E}_T \text{trkmet})$ large (it could point anywhere)

Missing energy from neutrinos trkMET & MET aligned



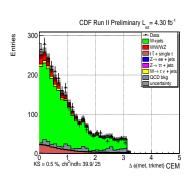
Checking the QCD II

- $\not\!\! E_T$ is normally calculated at calorimeter level
- Can also estimate it using tracks
- ullet True $ot\!\!\!E_T: \Delta\phi(
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- Fake \mathscr{E}_T : $\Delta \phi(\mathscr{E}_T \text{trkmet})$ large (it could point anywhere)



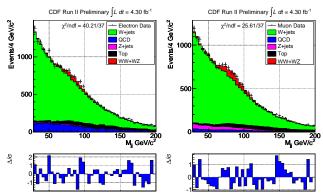
Checking the QCD II

- E_T is normally calculated at calorimeter level
- Can also estimate it using tracks
- $\qquad \text{True } \cancel{E}_T: \ \Delta\phi(\cancel{E}_T, \mathsf{trkmet}) \\ \mathsf{small}$
- Fake \mathbb{Z}_T : $\Delta \phi(\mathbb{Z}_T \text{trkmet})$ large (it could point anywhere)



Fit to M_{jj} distribution

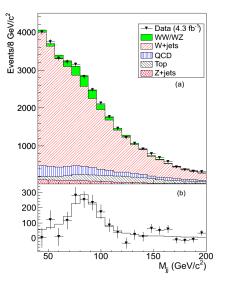
- Binned fit to the mjj shape taking as templates the histograms:
 - \bigcirc W+jet \longrightarrow completely free in the fit
 - SIGNAL (WW and WZ)
 - **QCD** \longrightarrow gaussian constraint to the value found in the \mathcal{E}_T fit with 25% width.
 - Top+single top : constrained to the measured cross section
 - 5 Z+jet: constrained to the measured cross section



Systematic summary

- Consider two classes of systematics:
 - systematics affecting signal extraction
 - systematics affecting signal cross-section
- Dominant Systematic for signal extraction:
 - **③** Shape of W+jets: evaluated changing the Q^2 of the MonteCarlo generator (6%).
 - ② Jet energy scale: evaluated varying the Jet energy scale up and down of 1 σ (6%).
 - QCD Shape: (4%).
- Dominant Systematic for signal cross-section:
 - Luminosity 6%
 - ② ISR/FSR 2%

Fit to data and results

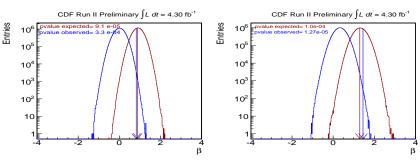


With $1582 \pm 275 ({\rm stat.}) \pm 107 ({\rm syst.})$ events this is the first observation (5.24σ) of $WW+WZ \rightarrow \ell \nu + {\rm jets.}$ The resulting x-sec is

$$\sigma(WW + WZ) = 18.1 \pm 3.3(\text{stat.}) \pm 2.5(\text{stat.})$$

that is in agreement with SM expectation (15 ± 0.9 pb).

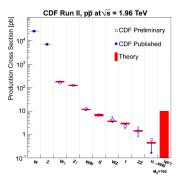
Significance estimation



- Generate one toy for each combination of the N_{syst} i.e. in each sample some of the systematic are varied.
- For each sample evaluate how many times the toy results exceeds the value observed in data and choose the worst
- The combined p-value is $8.56*10^{-08}$
- 5.24 σ found where 5.09σ was expected

Conclusion

- Measured the cross section of WW/WZ $\rightarrow l \nu$ +jets
 - Opens the way to diboson studies with jets
- PRL published on march 2010: Phys. Rev. Lett. 104, 101801 (2010)
- Webpage for the update

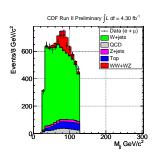


- Work in progress:
 - **①** Tagged analysis for $WZ o l
 u + ext{ bjets (} \sigma = 0.12 ext{ pb)}$
 - Model independent search for di-jet resonances associated with a W

Backup

Search for a di-jet resonance

- Using the same final state to look for resonances in the di-jet mass spectrum
- \bullet Since we are looking at heavier resonances moved the E_T cut to 30 GeV



General strategy for a bump hunt in the di-jet invariant mass:

- Mass search region: $120 < M_{ij} < 200 \text{ GeV/c}^2$
- The only assumption is that the signal width is compatible with the experimental resolution: so the expected width is

$$\sigma_{resonance} = \sigma_W \sqrt{\frac{M_{jj}}{M_W}}$$

Fitting procedure

- ullet Combined χ^2 fit to electrons and muons
- 6 templates:

 - QCD (constrained to its fraction with 25 % error)
 - \bigcirc Z + jets (constrained to the measured cross section)
 - top & single top (constrained to the measured cross section)
 - \bigcirc WW/WZ (constrained to the theoretical cross section)
 - onew resonance (gaussian with width related to the mass)
- Procedure:
 - **①** Fit the data without the resonance \rightarrow evaluate χ^2
 - ② Fit the data with the resonance ightarrow evaluate χ^2
 - **③** We add 3 degrees of freedom to the fit (mass, separate yields) so the $\Delta \chi^2$ should have the distribution of a χ^2 with 3 degrees of freedom.
 - **1** Verify the behaviour of the $\Delta \chi^2$ with toys with trial factor.